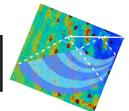
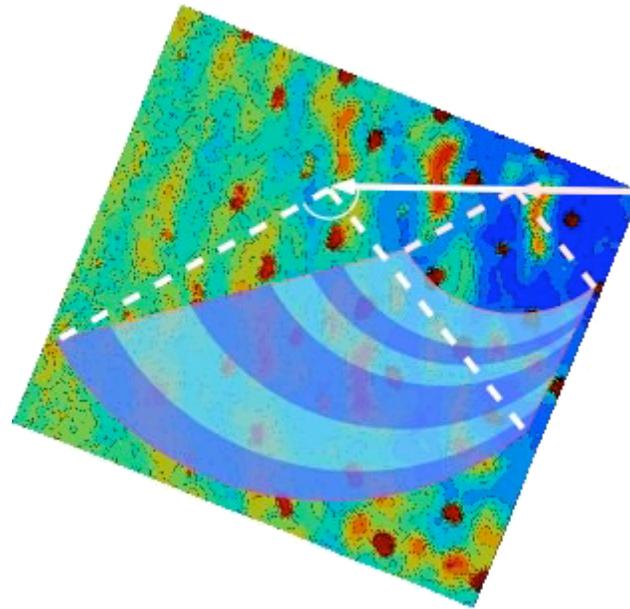




Cross Correlation Spectroscopy at Pulsed Neutron Sources

Stephan Rosenkranz and Ray Osborn

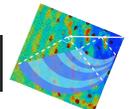




Goal

- Measure single crystal diffuse scattering $S(Q, \omega=0)$ at a pulsed neutron source
 - Over a large volume of reciprocal space
 - With energy discrimination
 - Good Q-resolution (at low Q's)
 - “large” samples (~cm's)
 - Low background

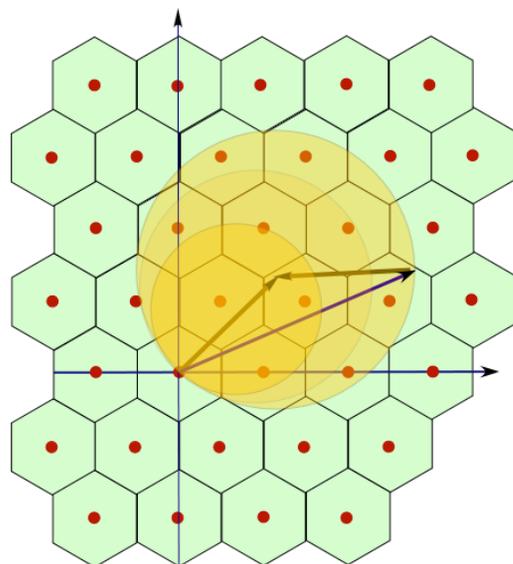
➔ Cross-correlation Spectrometer





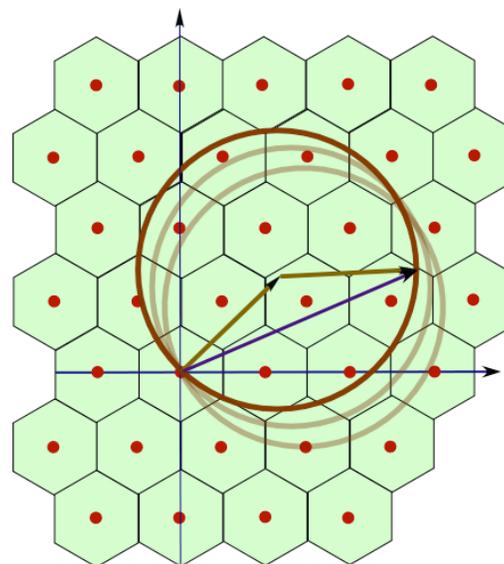
Conventional Techniques

White Beam

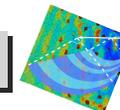


NO energy discrimination

Fixed k_i



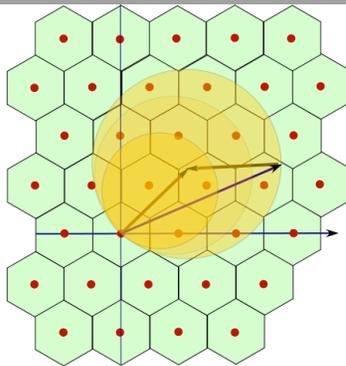
NOT efficient





Proposal

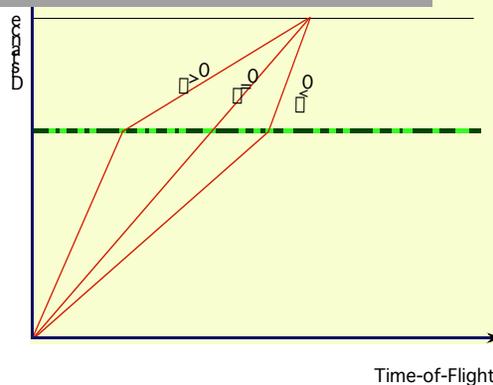
TOF Laue Diffractometer



“combine white beam and fixed k_i ”

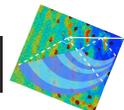
1. TOF Laue Diffractometer
 - highly efficient data collection
 - wide dynamic range in Q
2. Statistical Chopper
 - elastic energy discrimination
 - optimum use of white beam

Statistical Chopper



$$C(j,k) = \prod_{i=i_{\min}}^{i_{\min}+N} A(i \square k) S(i, j \square i) + B(j)$$

$$S(i, j) = \frac{2}{N+1} \prod_{k=1}^N A(i \square k) C(i + j, k) \square \frac{2}{N+1} B(i + j)$$





Principal of Correlation Technique

- First proposals and tests for high-efficiency time-of-flight methods ~ 1968
- Beam transmission modulated in time by $M(t)$ with $0 \leq M(t) \leq 1$
- Autocorrelation for pseudorandom sequence

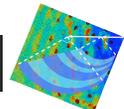
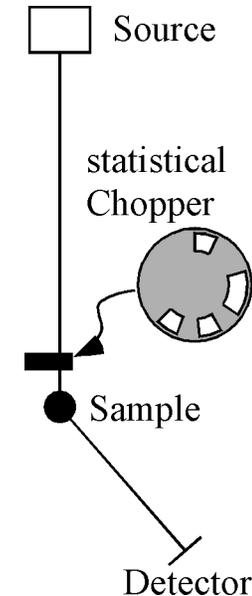
$$A^{MM}(\tau) \equiv \frac{1}{T_S} \int_0^{T_S} M(t)M(t-\tau)dt = c_1 \cdot \tau(\tau)$$

- Signal $I(t_0, t)$ due to sample scattering $S(t, t_1)$ and background $B(t)$

$$I(t_0, t) = \int_0^{T_S} M(\tau - t_0)S(t, \tau)d\tau + B(t)$$

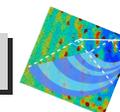
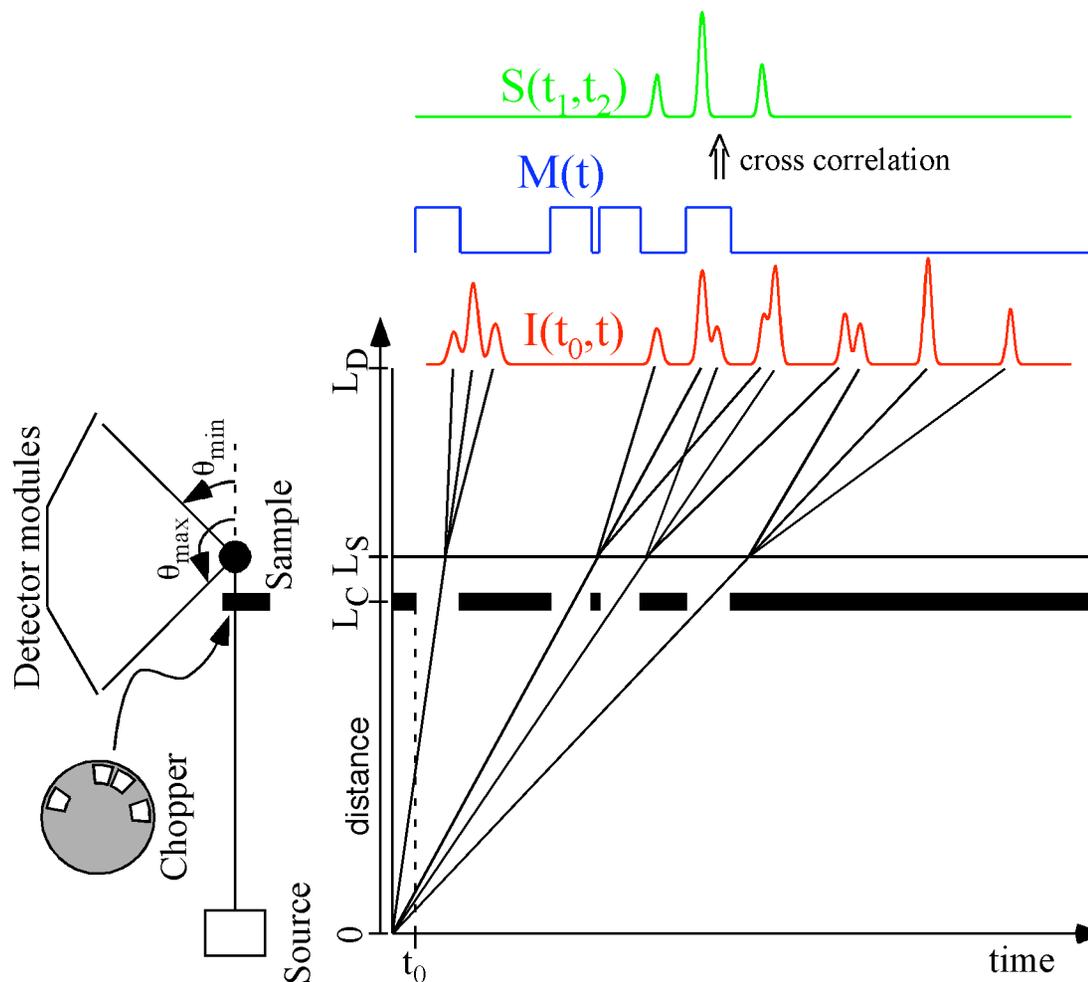
- Cross correlation

$$C(t, t_1) \equiv \frac{1}{T_S} \int_0^{T_S} M(t_1 - t_0)I(t_0, t)dt_0 = c_1 \cdot S(t, t_1) + c_2 \cdot B(t)$$



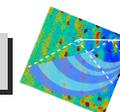
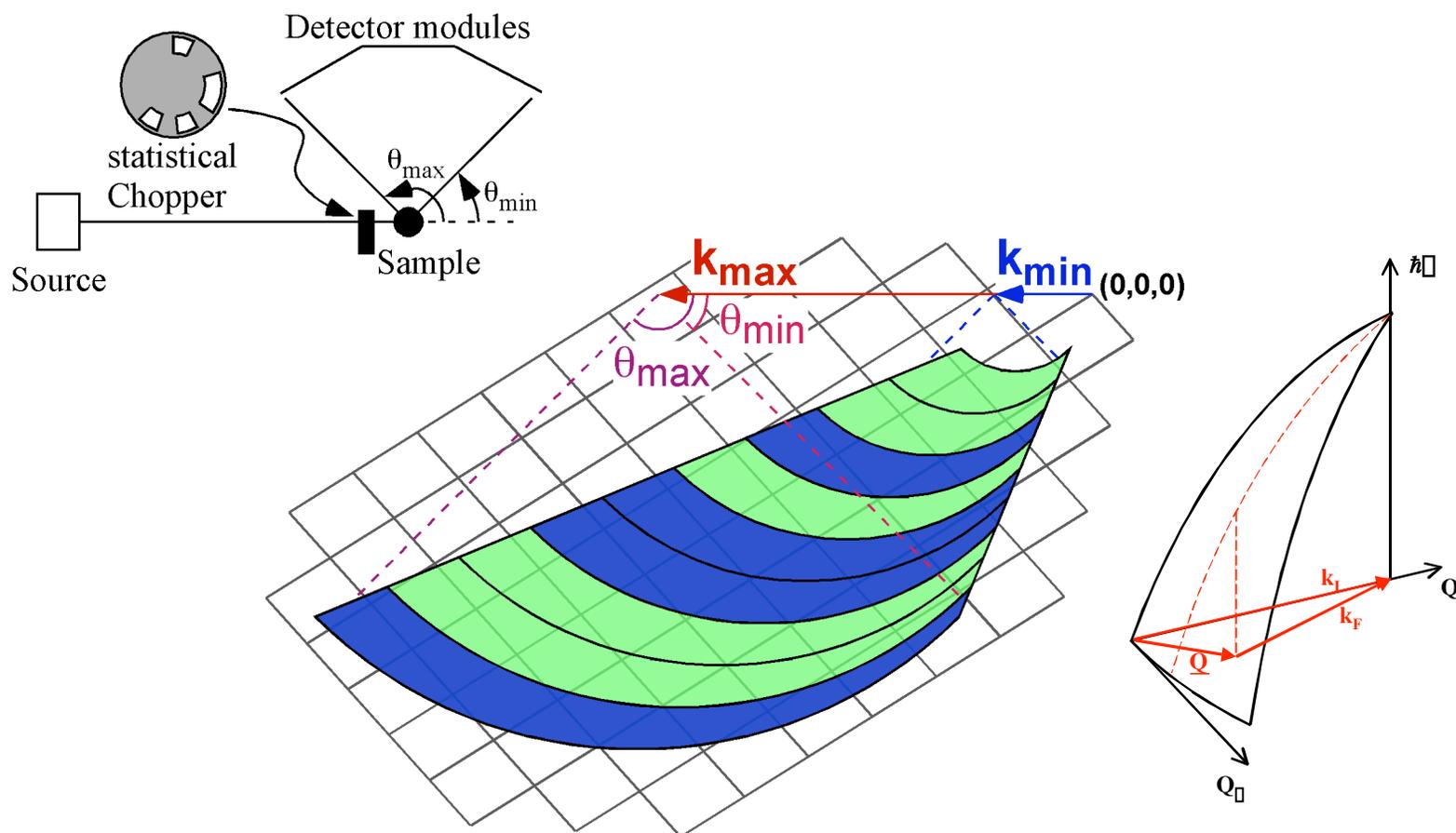


Cross Correlation t-D diagram





Elastic coverage

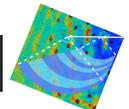




Statistics

- time-channel notation
$$C(j,k) = \frac{1}{N} \sum_{l=1}^N M(j-l)I(l,j-k)$$
- Relative signals
$$S(i,j) = \frac{2}{N+1} \sum_{k=1}^N A(i-k)C(i+j,k) - \frac{2}{N+1} B(i+j)$$

$$\sigma_{ij} = \frac{S(i,j)}{\bar{S}}, \quad \sigma_j = \frac{B(j)}{\bar{S}}$$
- Gain Factor
$$G = \frac{[\text{rel. var. } S(i,j)]_{\text{conv}}}{[\text{rel. var. } S(i,j)]_{\text{corr}}} \approx \frac{m(1-c)}{1-2c + \frac{2c+cN}{\sigma_{ij} + \sigma}}$$
- Average gain Factor
$$G_{\text{av}} \approx (1-c) \frac{1+\sigma}{1 + \frac{\sigma}{m}}$$



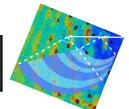


Does it work?

- Test performed on LRMECS at IPNS
- Chopper with 251 elements running at ~ 200Hz, asynchronously with the source (30Hz)
- Modified data acquisition module
- Measured standard vanadium sample for 4.5h
(Test's terminated because of accelerator breakdown)

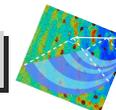
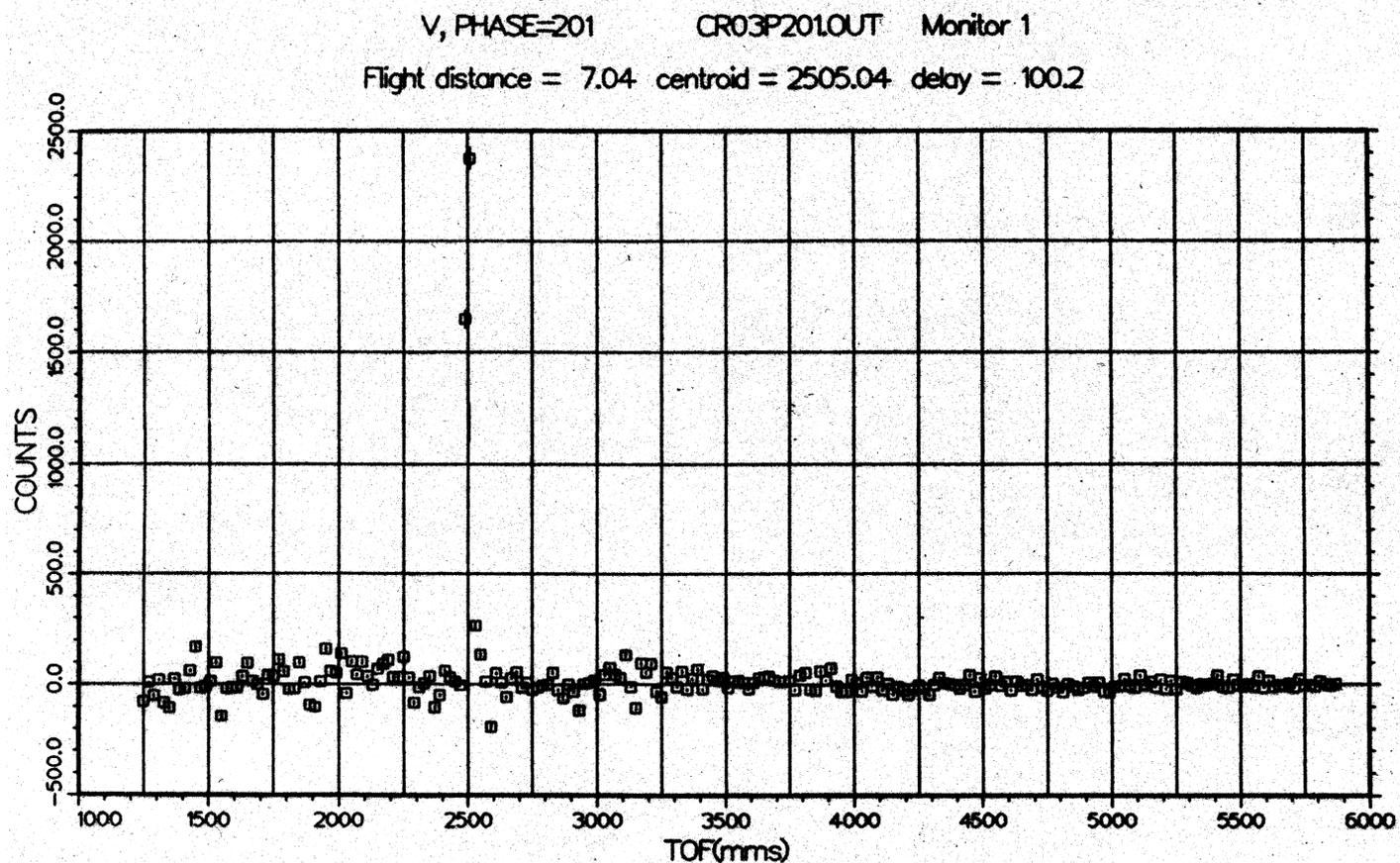
TEST OF A CORRELATION CHOPPER AT A PULSED
SPALLATION NEUTRON SOURCE*

R. K. Crawford,** J. R. Haumann,⁺ G. E. Ostrowski,**
D. L. Price,⁺⁺ and K. Sköld[†]
Argonne National Laboratory
Argonne, Illinois 60439



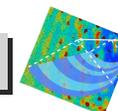
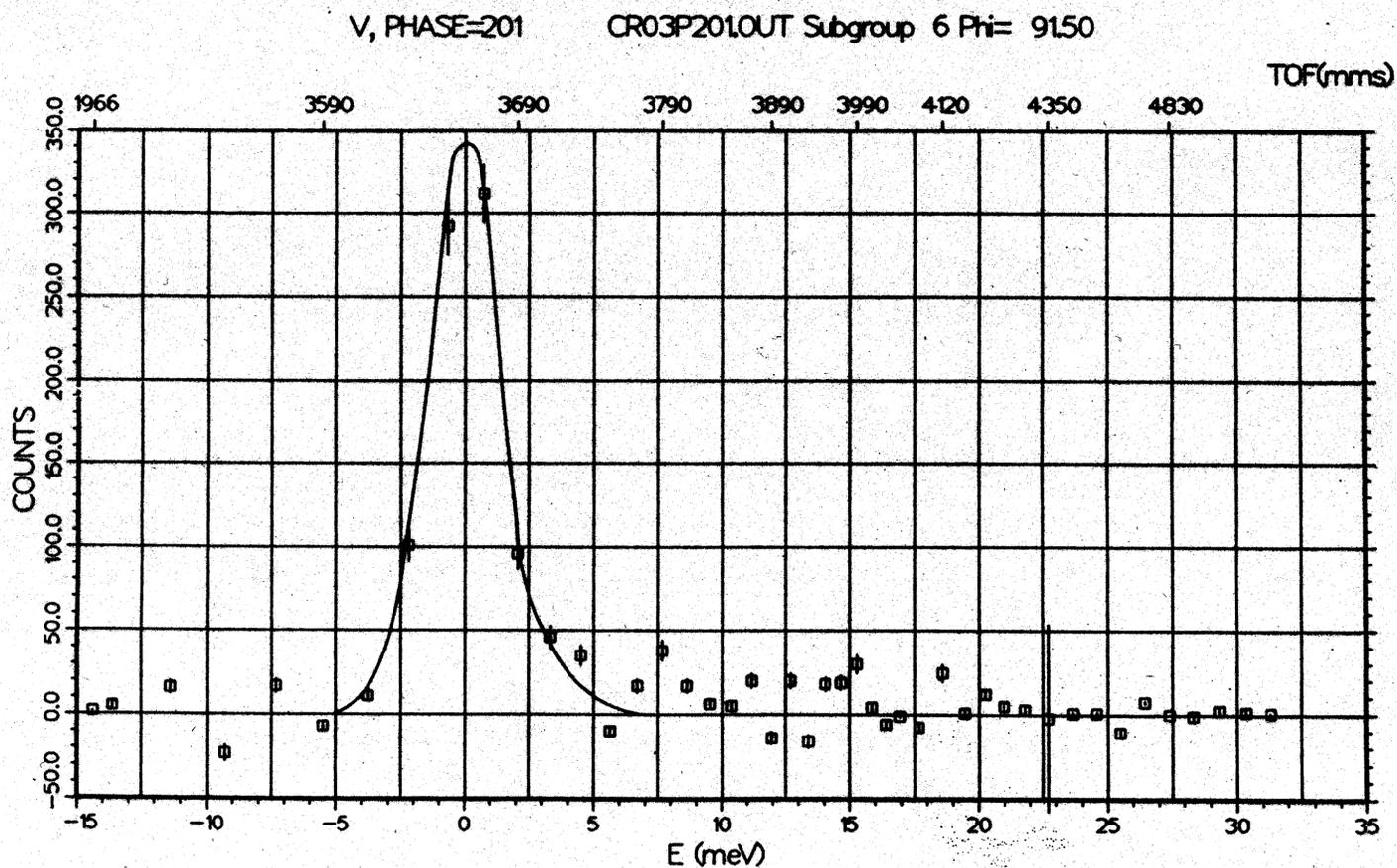


Monitor





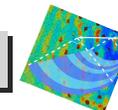
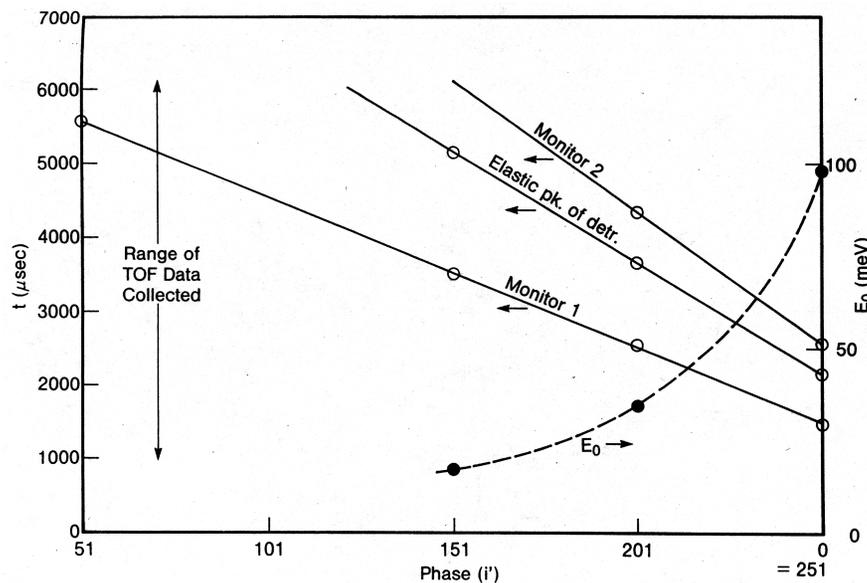
Vanadium





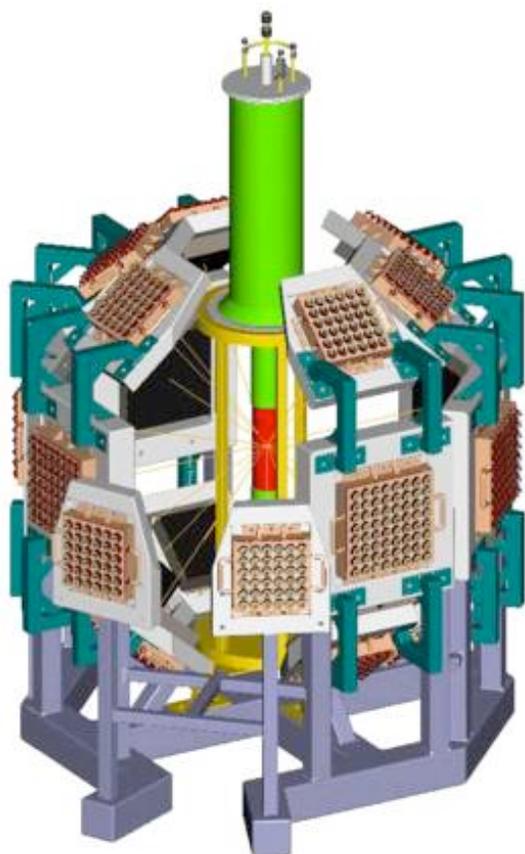
Elastic peaks

- Within accuracy of experiment, source pulse at $t=0$ for all neutrons
- Energy resolutions = 1.5 ($E_0=17.5$), 3.2 (25), 16.2 meV FWHM, slightly worse than typical value for Fermi choppers

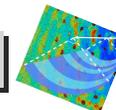




Option for existing Instruments?

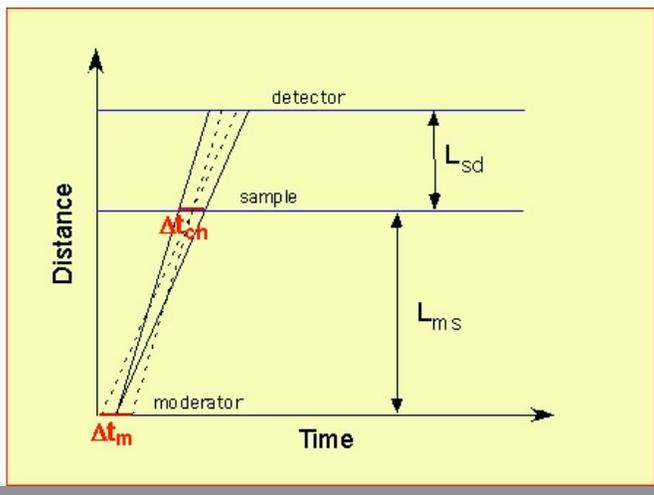


- Diffuse scattering makes particularly stringent demands on instrumental backgrounds
 - The scattering is very weak
 - It is broadly distributed in Q-space
 - It is not easily identified
e.g., by Brillouin zone symmetries
- The current design of SCD envisages an open instrument
 - Bragg peaks are relatively easy to identify and integrate
 - However, there will be backgrounds from:
 - Air scattering
 - Extra sample environment windows
- We believe that it will be essential to have an evacuated secondary flight path





Resolution



Q-resolution

SCD plans to have reasonable resolution

$$\Delta d/d \sim 0.01$$

with $L_{ms} \sim 15\text{m}$ and $L_{sd} \sim 1\text{m}$

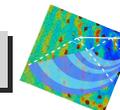
There are two problems with this value:

1. It is calculated at 90° .
 - We need good resolution at low Q
2. It neglects sample size.
 - We must accommodate large samples.

ΔE -resolution

- The secondary flight path of 1m is too small to give adequate energy resolution.

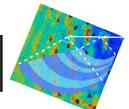
$$\frac{\Delta E}{E_i} = \sqrt{\left(\frac{\Delta t_{ch}}{t_{ch}}\right)^2 + \left(\frac{L_{ms}}{L_{sd}}\right)^2 + \left(\frac{\Delta t_m}{t_{ch}}\right)^2}$$





Summary

- It is unlikely that the demands of the crystallography and diffuse scattering communities can be satisfied by a single diffractometer. Diffuse scattering requires:
 - **low backgrounds**
 - **good resolution at low-Q**
 - **a longer secondary flight path for adequate energy resolution**
- Furthermore, the potential scientific demand for diffuse scattering is sufficient to justify a dedicated diffractometer, since:
 - **it overcomes previous technical limitations**
 - **there is increased interest in the science of complex disorder**
- We plan to hold a workshop NOW to strengthen the scientific case for such an instrument and establish the relevant scientific community.





Future developments

- Thorough analysis of cross-correlation technique
 - Optimized statistical sequences required for different kinds of diffuse scattering
 - Best method for generating and phasing the pseudo-random pulses to the source
 - Most practical wavelength range
 - Design parameters for an optimized chopper



- Build prototype Chopper and test performance for measuring single crystal diffuse scattering



We just received funding from DOE and IPNS

