

Emergent Excitations in Geometrically Frustrated Magnets

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- ❑ Issues on geometrical frustrated insulating and metallic magnets
 - What is the nature of the spin liquid phase?
 - Any novel phase transition?
 - What is the effect of GF in the metallic systems?
- ❑ Neutron scattering studies on insulating and metallic spinels AB_2O_4
 - $ZnCr_2O_4$ and $Li_{1-x}Zn_xV_2O_4$
- ❑ Future studies...

Collaborators:

Experiments:

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S. Rosenkranz, R. Osborn (ANL)

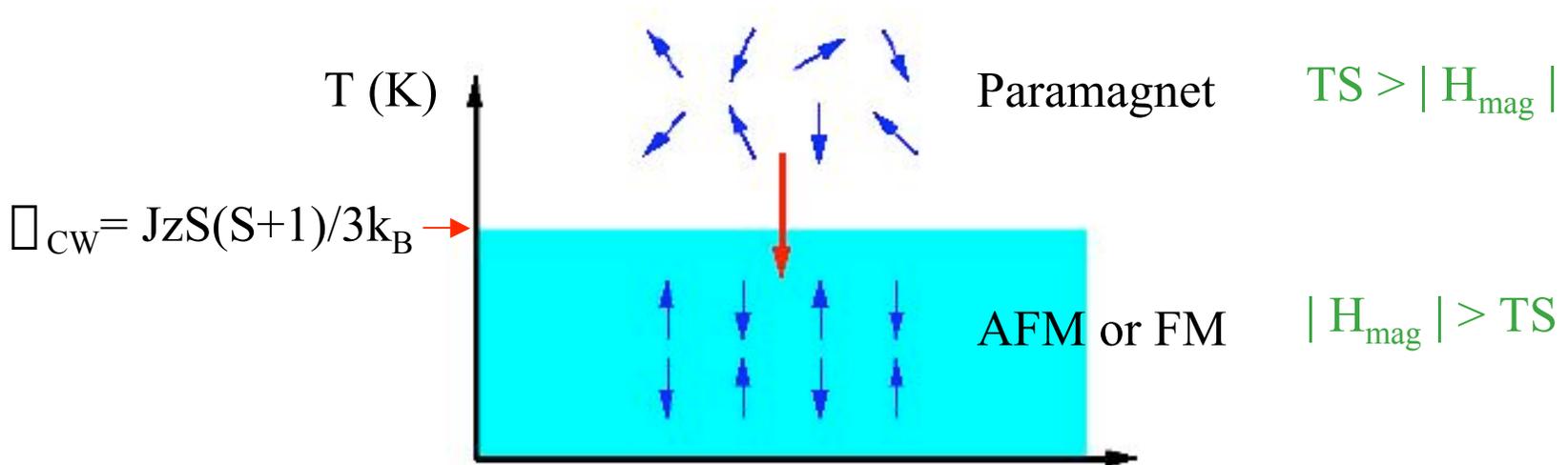
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Theory:

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Phase Diagram for an Ordinary Magnet

$$F = H_{\text{mag}} - TS = - \sum J S_i S_j - TS$$

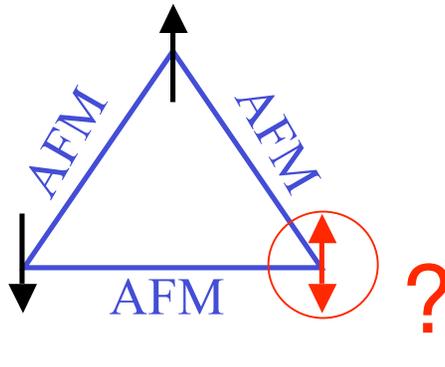


Ground State : Long range (anti)ferromagnetically ordered state.

Low Energy Excitations : linear spin waves around the ordered state.

Geometrical Frustration

A simplest example: a Triangle of three antiferromagnetic Ising spins

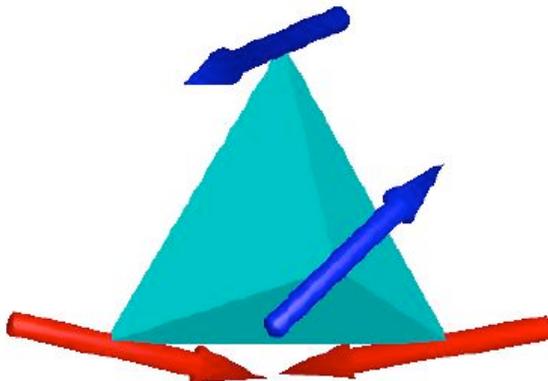


$$H = -J \sum S_i \cdot S_j$$

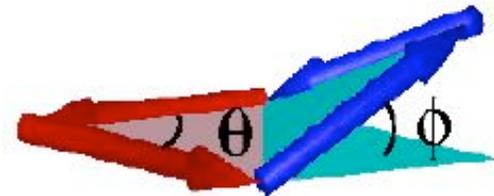
All exchange interactions can not be satisfied.

$$E = -|J|$$

A tetrahedron with four Heisenberg spins



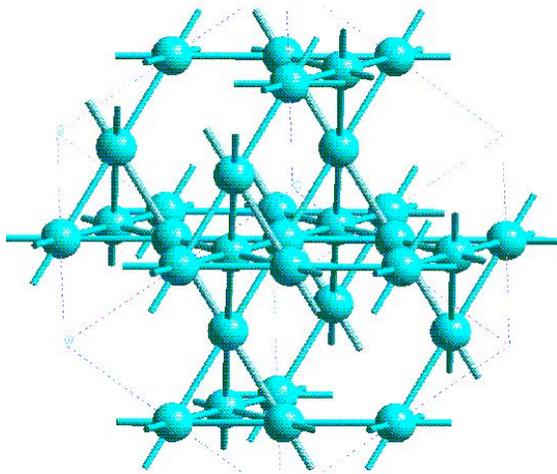
$$\sum S_i = 0$$



Zero energy modes
in the ground state manifold

Geometrical frustration leads to a large degeneracy in the ground state

Theory of spins with AFM interactions on corner-sharing tetrahedra



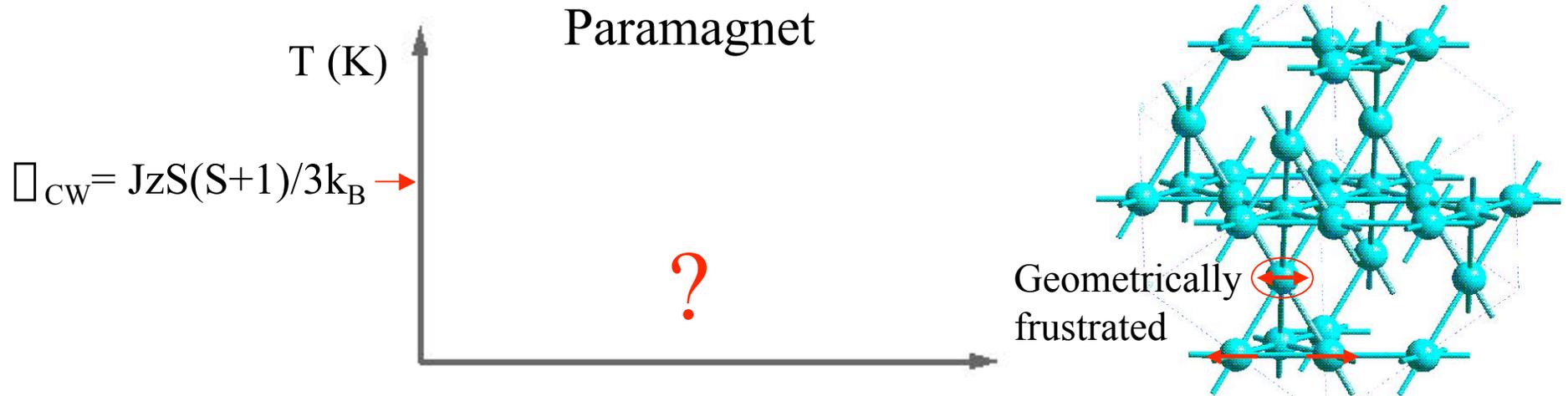
The simplest and most frustrated case is a system with isotropic uniform nearest neighbor interactions only:

$$H = -J \sum S_i \cdot S_j$$

SPIN TYPE	SPIN Value	LOW T PHASE	METHOD	REFERENCE
Isotropic	$S=1/2$	Spin Liquid	Exact Diag.	Canals and Lacroix PRL '98
Isotropic	$S=\infty$	Spin Liquid	MC sim.	Reimers PRB '92 Moessner, Chalker PRL '98

Phase Diagram for GF Insulating Magnets

$$F = H_{\text{mag}} - TS = - \sum J S_i S_j - TS$$



➤ Nature of the Spin Liquid State

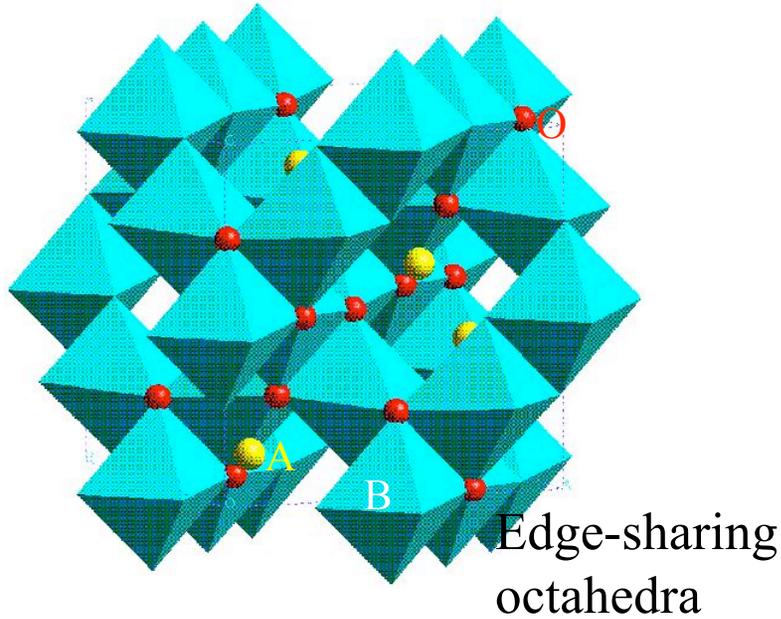
How are the fluctuating spins in the SL phase correlated with each other?
Zero energy mode?

➤ Novel Phase Transition?

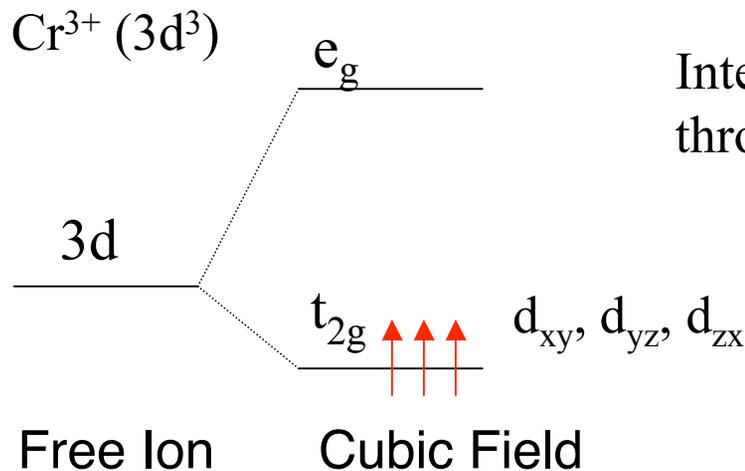
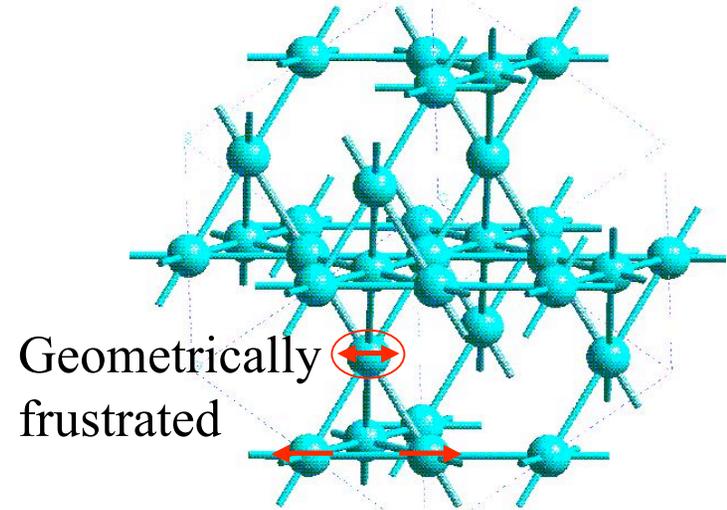
How does nature respond to the ground state degeneracy?

Why ZnCr_2O_4 ?

Space group $\text{Fd}\bar{3}\text{m}$



Lattice of B sites
: Corner-sharing tetrahedra

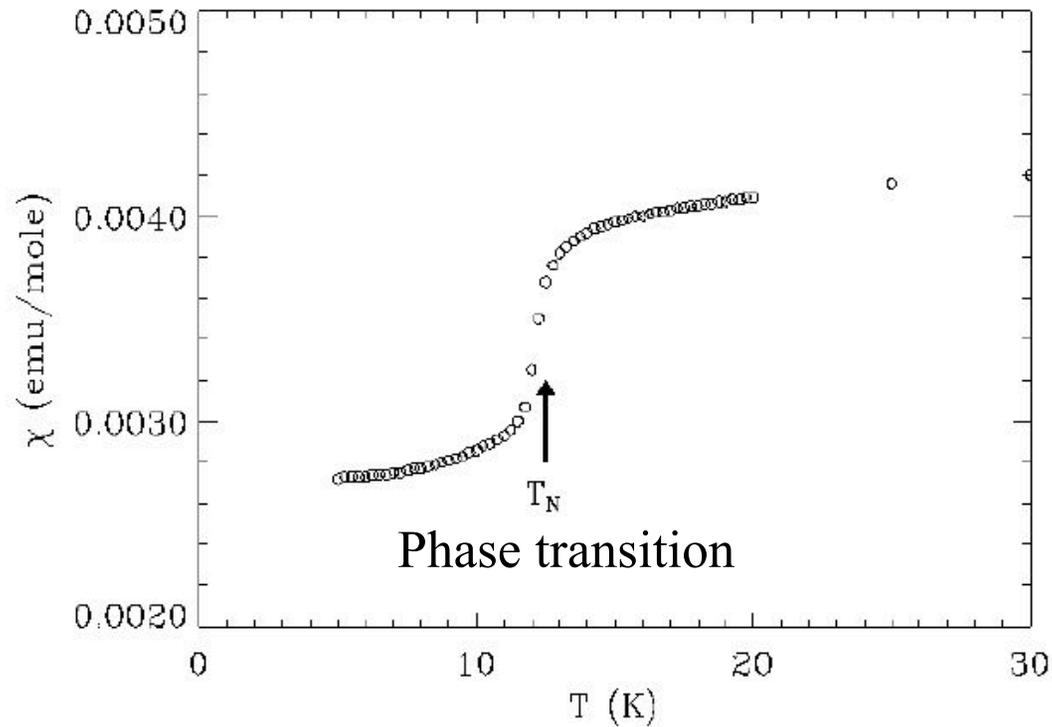


Interaction between nearest neighbor Cr^{3+} ions through direct overlap of t_{2g} orbitals is dominant.

Best candidate for a system with

$$H = -J \sum_{nn} \mathbf{S}_i \cdot \mathbf{S}_j$$

Magnetic Phase Transition in ZnCr_2O_4

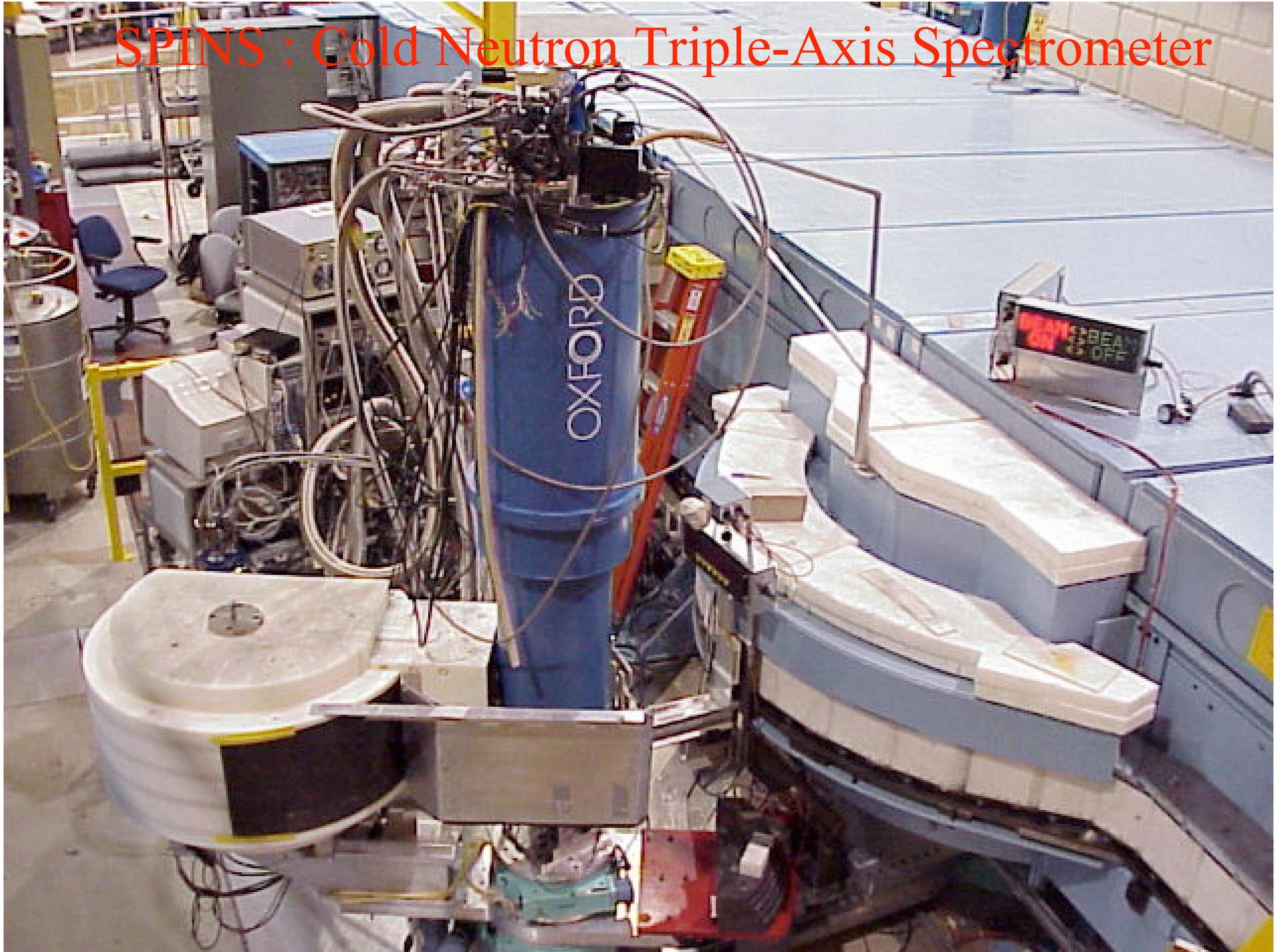


$$\square_{\text{CW}} = -390 \text{ K}$$

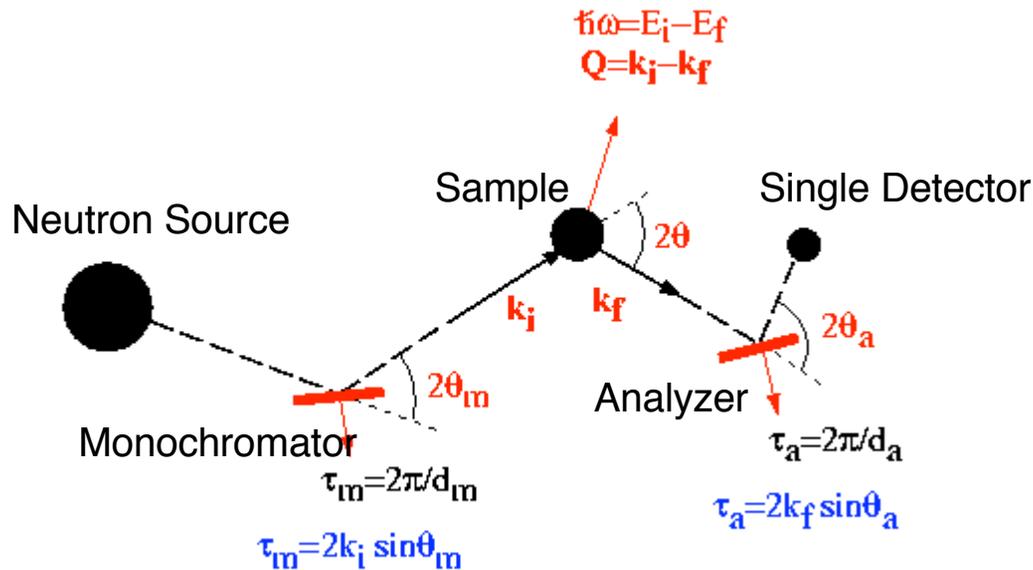
$$T_{\text{N}} = 12.5 \text{ K}$$

What is nature of the phase transition ?

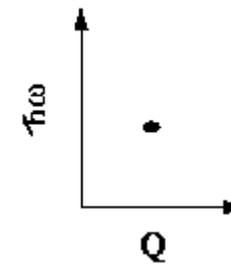
SPINS : Cold Neutron Triple-Axis Spectrometer



Conventional Triple-Axis Spectroscopy (TAS)



A single point at a time



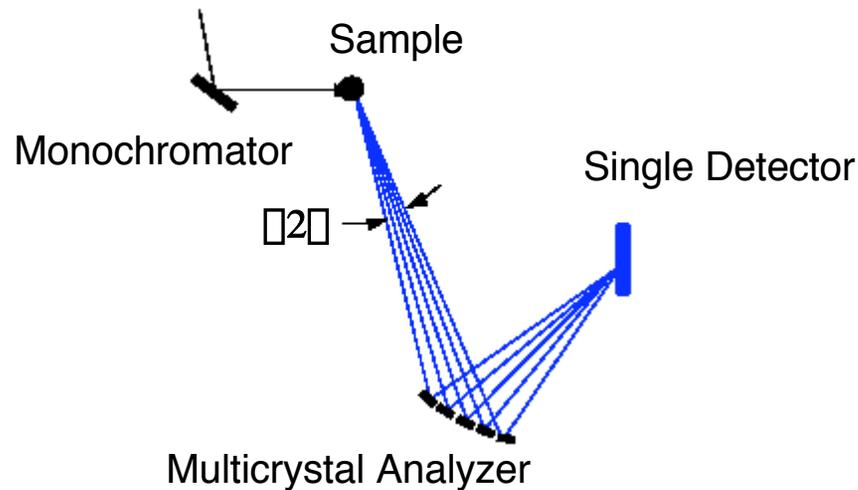
TAS is ideally suited for probing small regions of phase space

Shortcoming: Low data collection rate

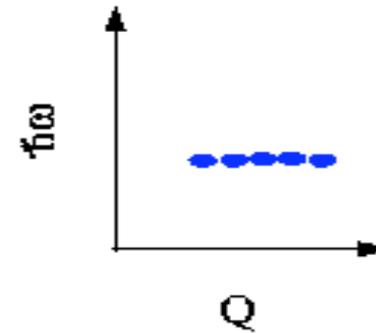
↓ **Improvement**

Multicrystal analyzer and position-sensitive detector

Horizontally Focusing (HF) Analyzer Mode



Relaxed Q-resolution



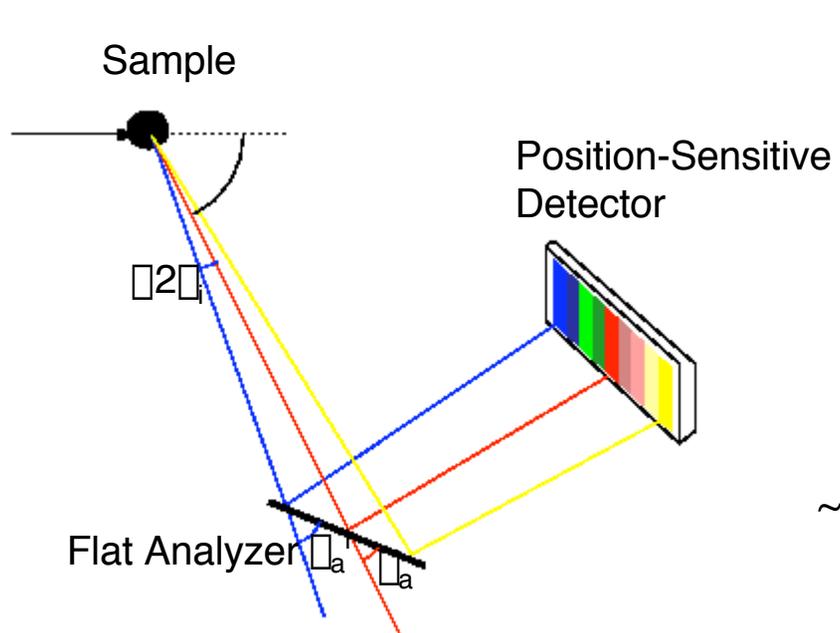
L = distance from sample to HF analyzer

w_a = total width of HF analyzer

$$2\theta = w_a \sin\theta_a / L \sim 9 \text{ degree for } E_f = 5 \text{ meV at SPINS}$$

Useful for studying systems with short-range correlations

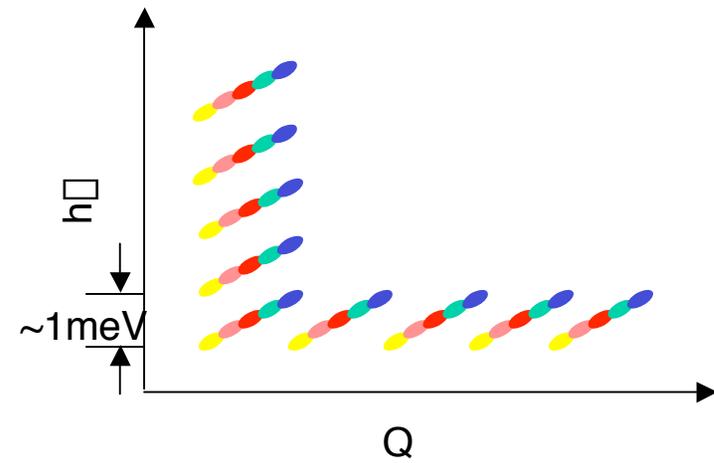
Multiplexing Detection System for TAS



$$\theta_a^i = \theta_a + \theta_i = \theta_a - \text{atan}(x \sin \theta_a / (L + x \cos \theta_a))$$

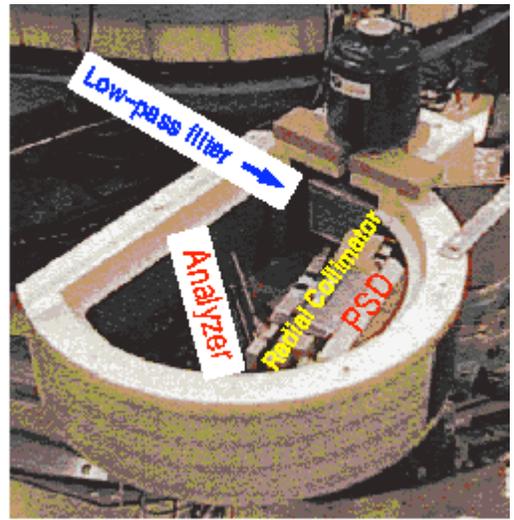
$$k_f^i = \theta_a / 2 \sin \theta_a^i$$

$$Q_i = k_i - k_f^i$$

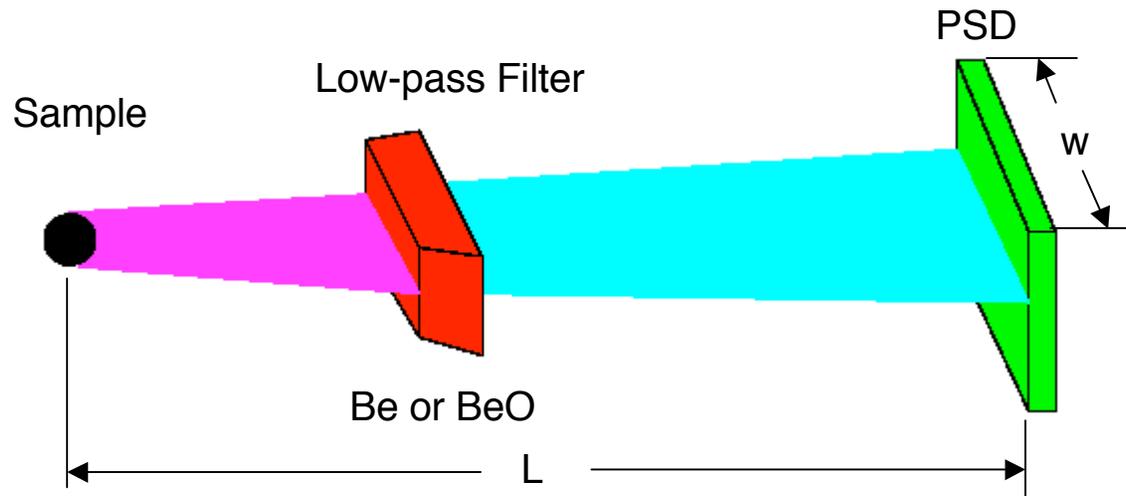


Probes scattering events at different energy and momentum transfers simultaneously

Survey ($h\nu$ -Q) space by changing the incident energy and scattering angle



Position-Sensitive Detector in Two-Axis Mode

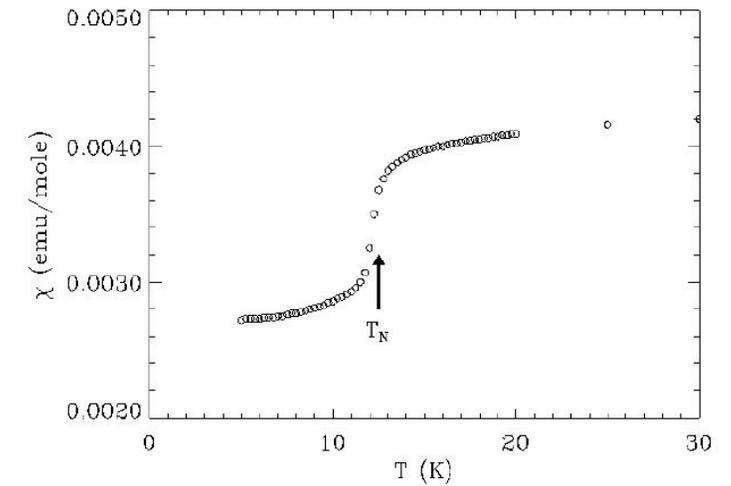
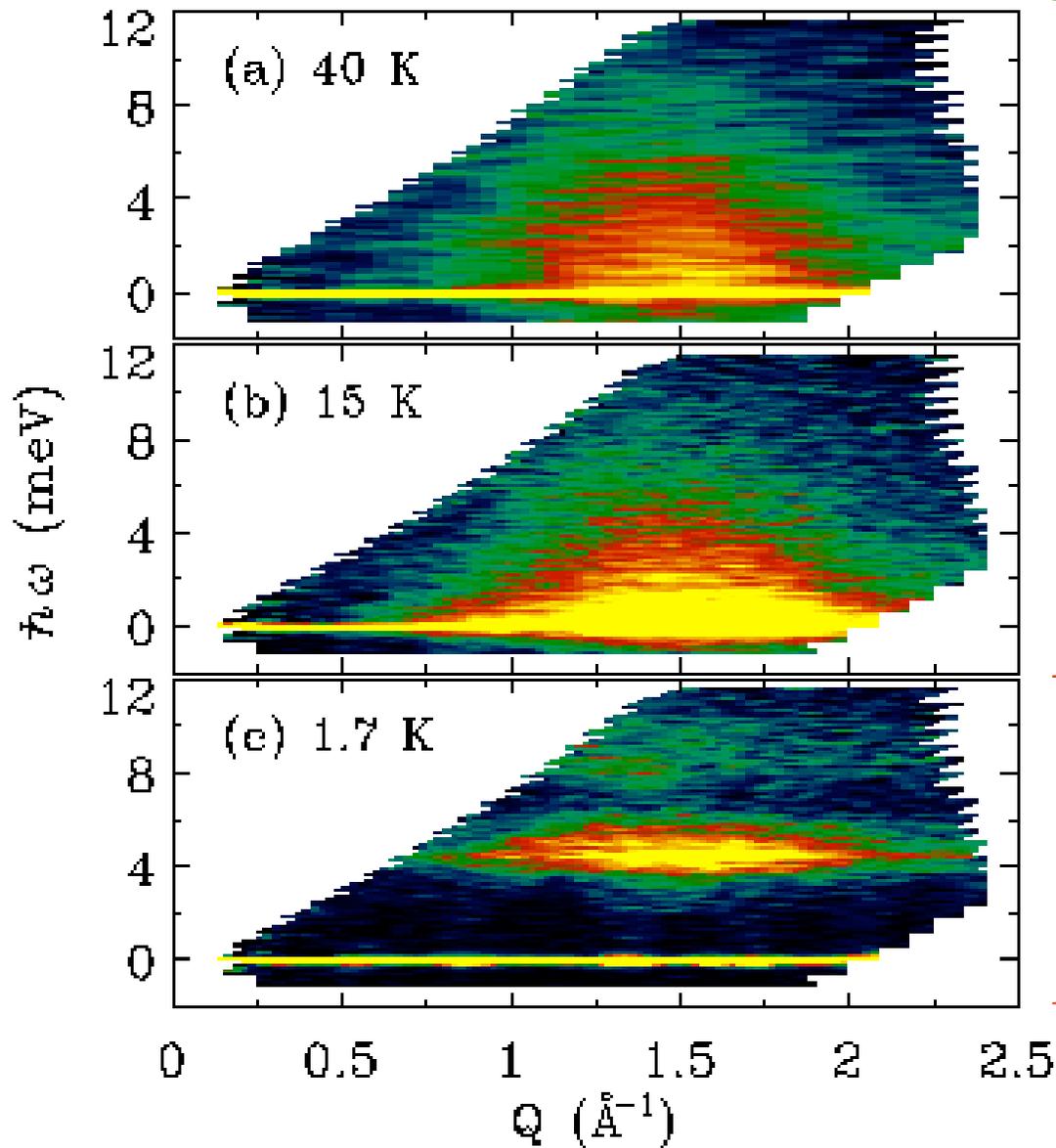


The filter passes only those neutrons with $0 < E_f < E_{\text{cutoff}}^{\text{filter}}$

The PSD measures $\int_{E_i - E_{\text{cutoff}}}^{E_i} S(Q, \square) d\square$

Large angular acceptance = $w / L \sim 11^\circ$ for SPINS

Neutron Scattering from ZnCr_2O_4



Cooperative Paramagnetic
fluctuations in
frustrated AFM

Local spin resonance
in magneto-elastic
LRO phase

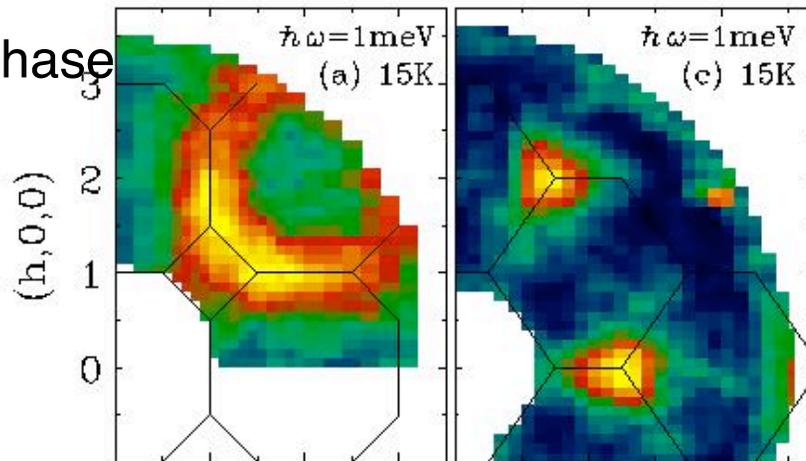
Nature of the Spin Liquid State of GF magnets

Emergent Excitations

Structure Factor of Spin Fluctuations in ZnCr_2O_4

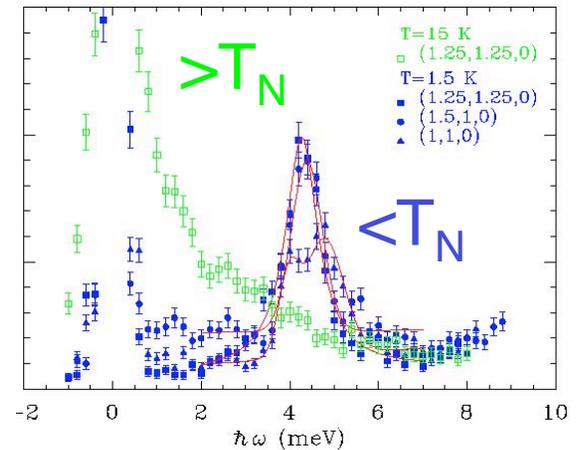
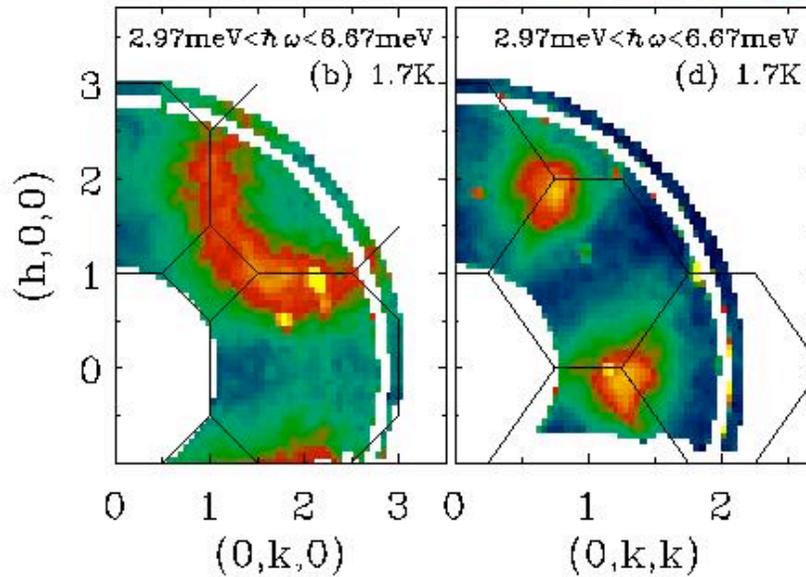
Spin liquid phase

$T > T_N$



Neel phase

$T < T_N$



Identical $S(Q)$

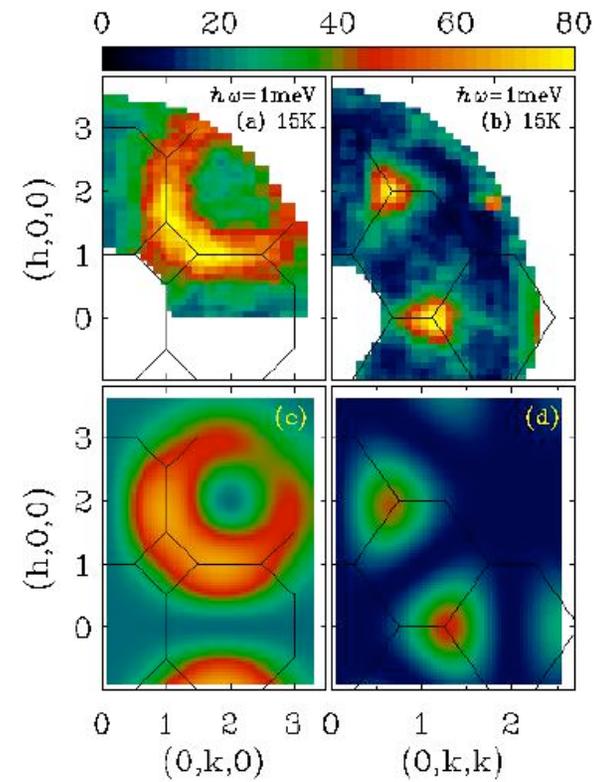
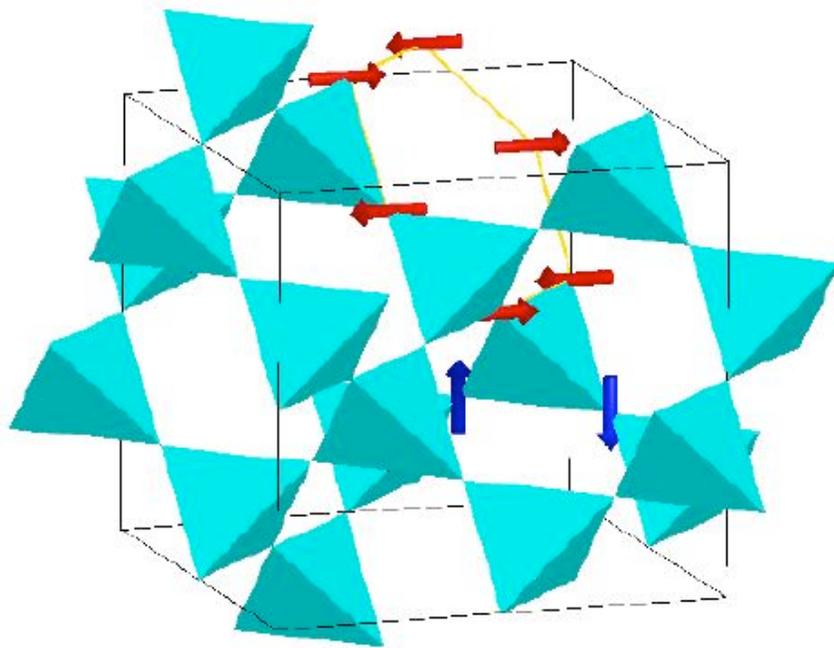
ZnCr_2O_4 single crystals



Extended Brillouin zone boundary scattering

Emergent Excitations in ZnCr_2O_4

SHL et al., Nature **418**, 856 (2002)



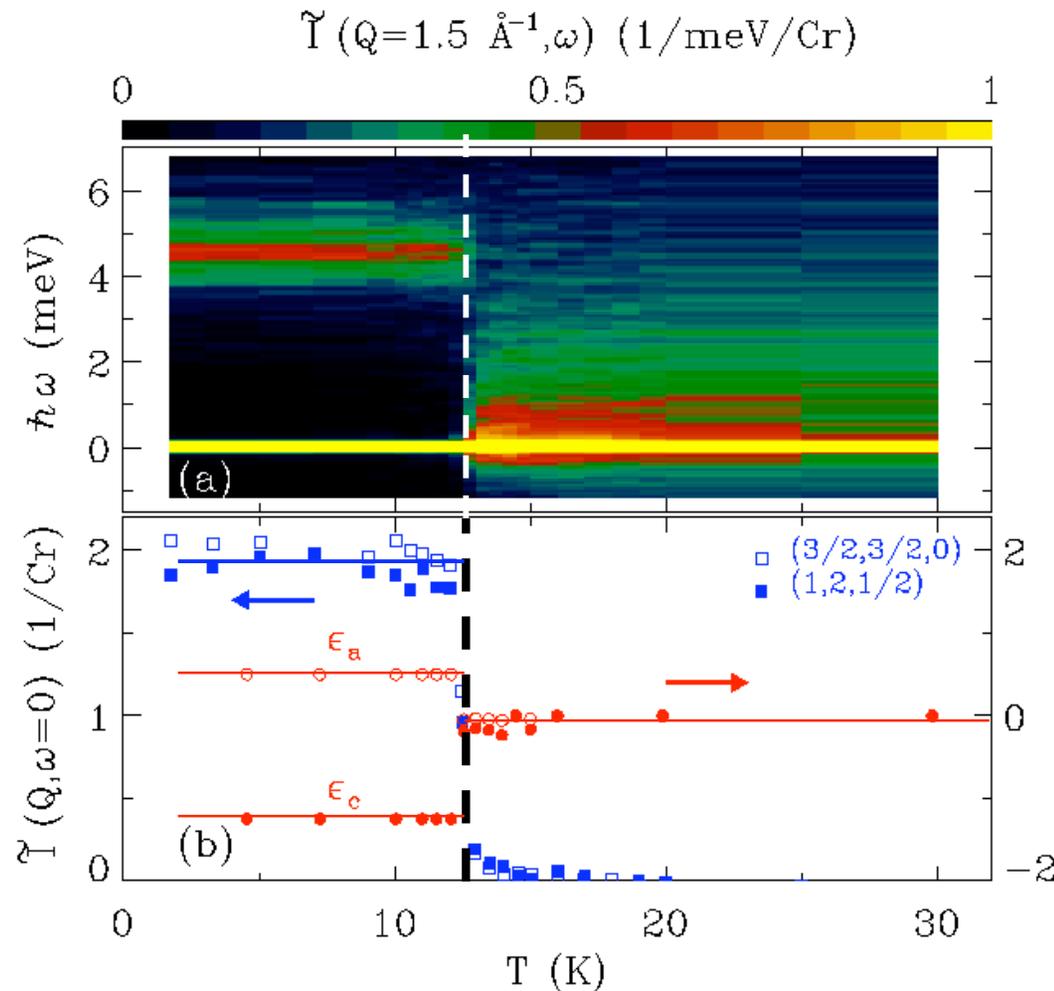
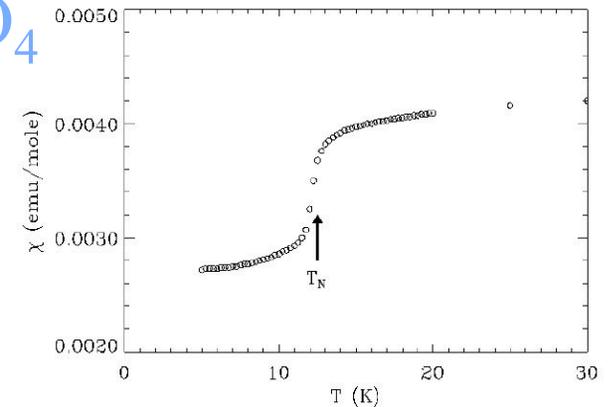
Hexagonal antiferromagnetic loops !

How does nature respond to the ground state degeneracy?

Spin-Peierls-like phase transition

Magneto-elastic phase transition in ZnCr_2O_4

How does nature respond to the ground state degeneracy?



Dynamics:

- Low energy paramagnetic fluctuations move to a resonance at 4.5 meV

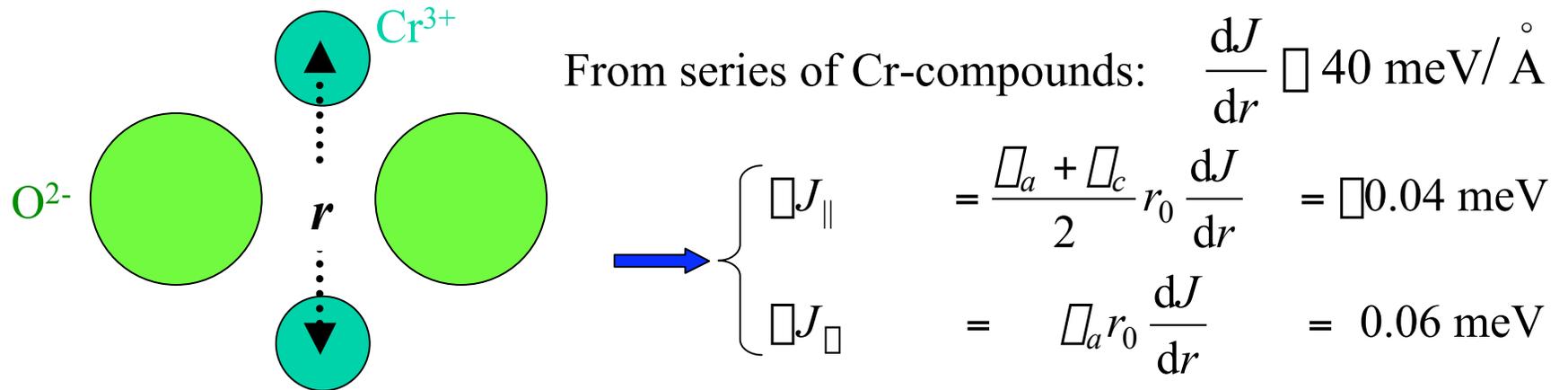
Statics:

- Staggered magnetization
- tetragonal lattice distortion

SHL et al., PRL 84, 3718 (2000)

Why does tetragonal strain encourage Neel order ?

Edge sharing n-n exchange in ZnCr_2O_4 depends strongly on Cr-Cr distance, r :



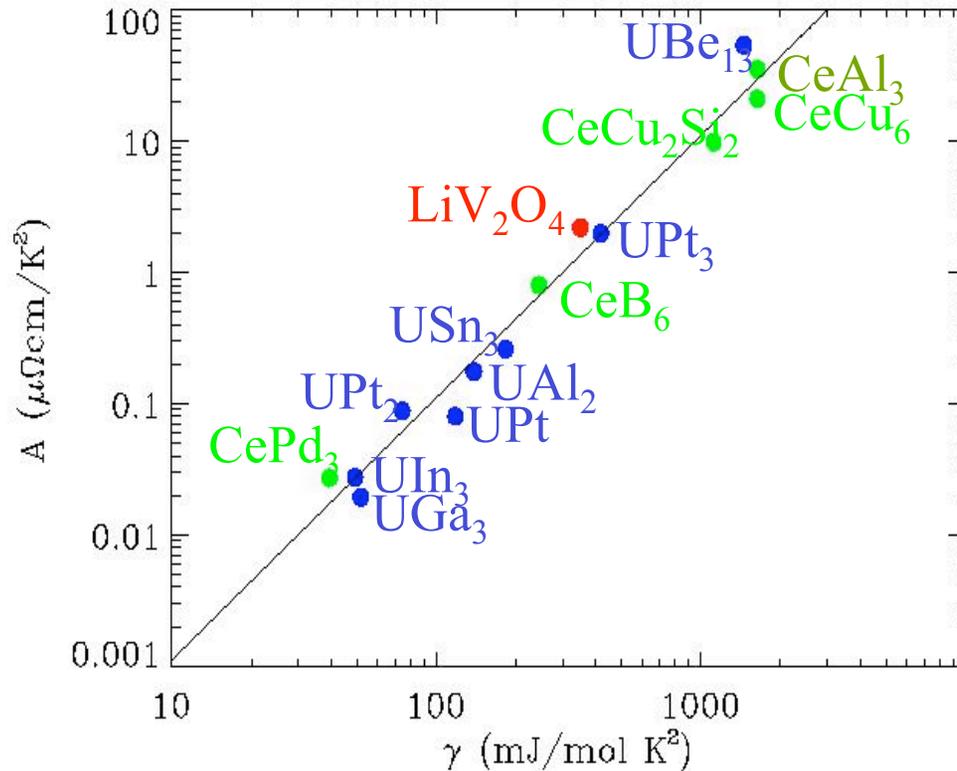
The effect for a single tetrahedron is to make 4 bonds more AFM and two bonds are less AFM. This relieves frustration!



**What is the effect of geometrical frustration
in a metallic system?**

LiV₂O₄: d-Electron Heavy Fermion

S. Kondo et al., PRL (1998)



Bulk measurement data from LiV₂O₄ at low T exhibit **Fermi liquid** behaviors

$$C_v \sim \gamma T$$

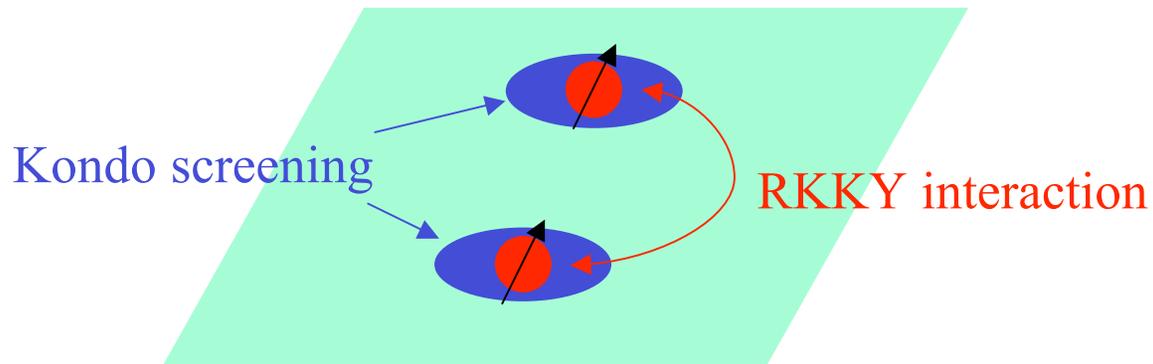
$$\rho \sim \rho_0 + AT^2$$

$$\chi \sim \text{const}$$

LiV₂O₄ with d-electrons is as heavy as UPt₃ !

Heavy Fermion

Heavy fermion behaviors with a heavy mass $m \sim 100-1000 m_e$ are usually found in **Ce- or U-based compounds** that have **two different types of electrons**: (1) **localized f-electrons** and (2) **conduction (s,p)-electrons**.

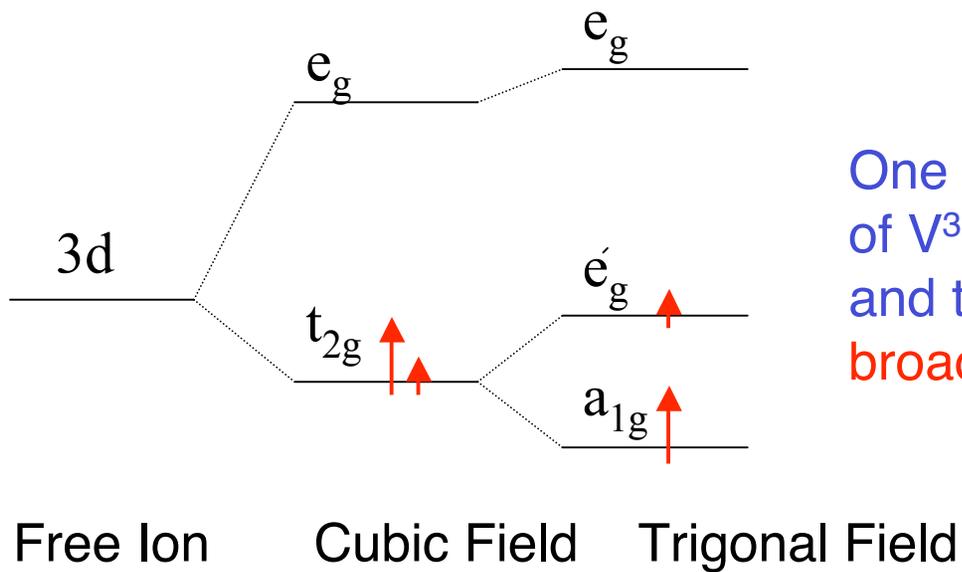


Issues : What is the mechanism for the HF behaviors in LiV_2O_4 ?
Why is LiV_2O_4 as heavy as UPt_3 ?
Is the HF state of LiV_2O_4 different from those in other Ce- and U-based HF compounds?

Mechanism for HF behavior in LiV_2O_4

1. Kondo screening

V.I. Anisimov et al., PRL 83, 364 (1999)

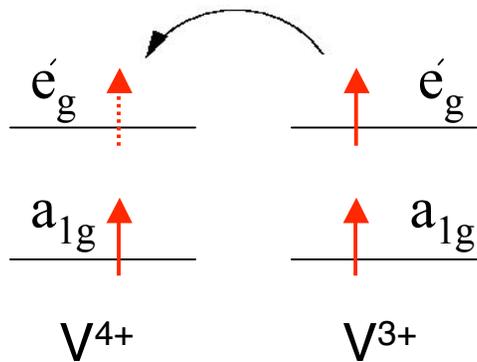


One electron of the $d^{1.5}$ configuration of $\text{V}^{3.5+}$ ion is in the localized a_{1g} orbital and the rest partially fills the relatively broad conduction e_g band

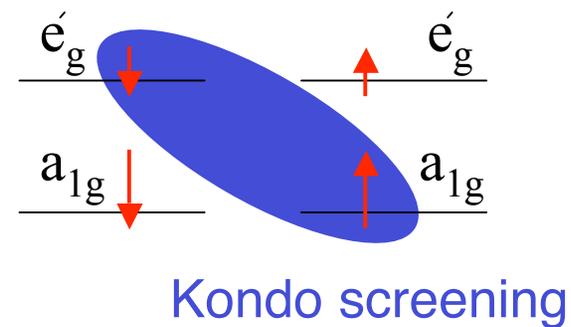
A Kondo screening possible !?!

Two competing interactions between neighboring local V moments:

FM double exchange interaction



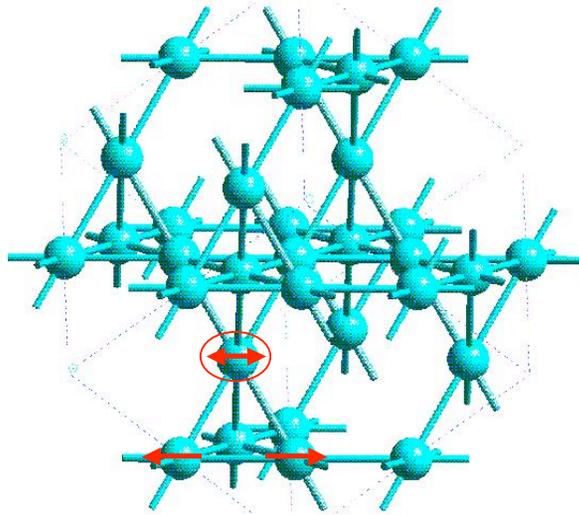
Kondo-induced **AFM** interaction



Mechanism for HF behavior in LiV_2O_4

2. Geometrical Frustration

V. Eyert et al., Europhy. Letters, 46, 762 (1999)

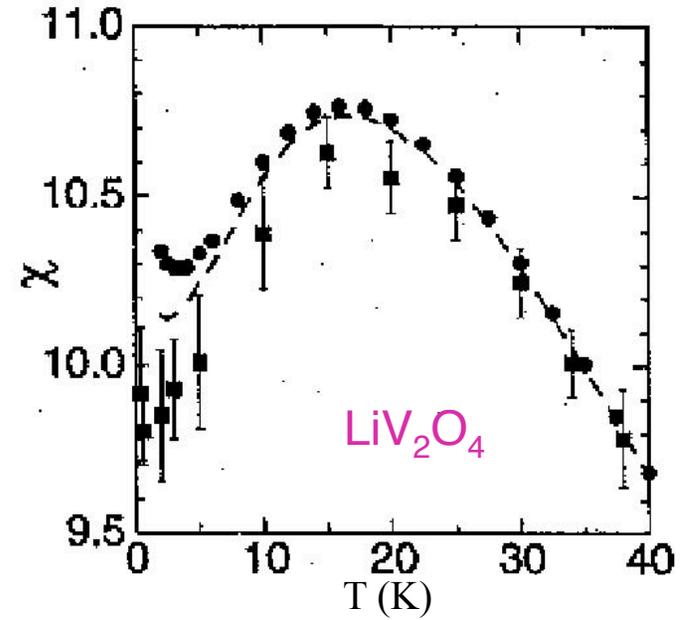
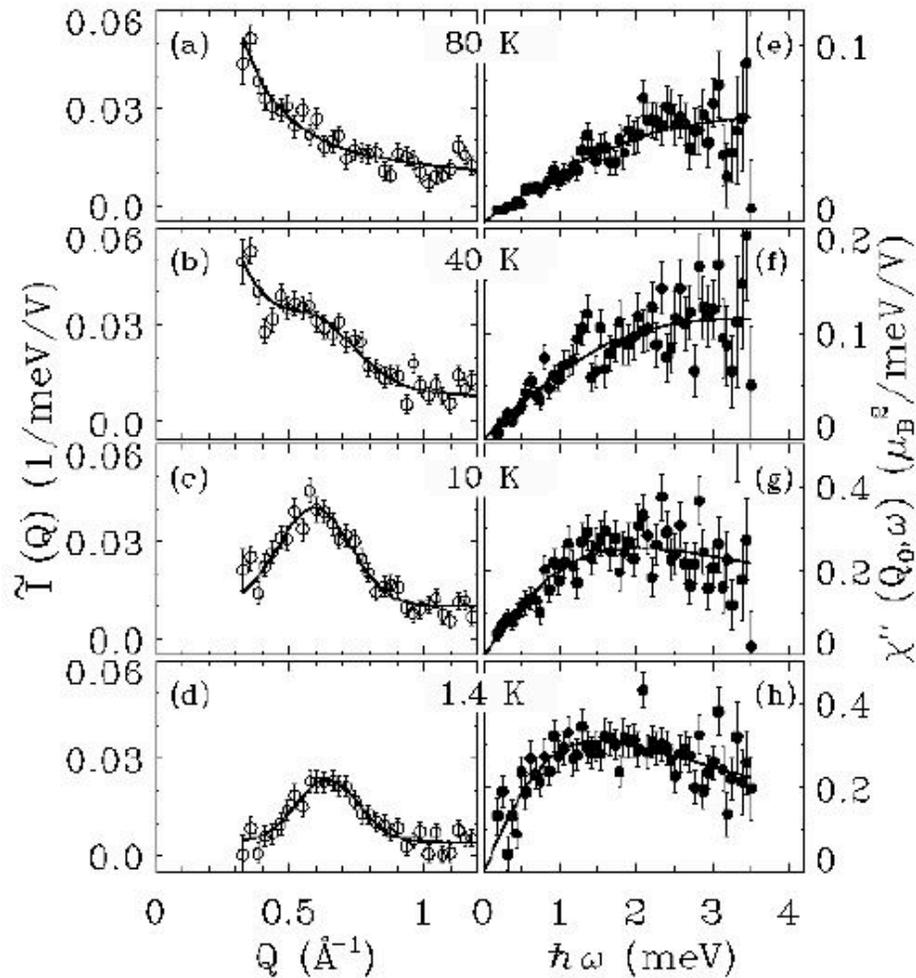


V moments interact **antiferromagnetically** in the frustrating lattice.

$$C_v = dE/dT = d(TS)/dT \sim S \text{ (Entropy)}$$

Geometrical frustration \rightarrow **Strong spin fluctuations**
 \rightarrow **Enhancement of specific heat (HF behavior)**

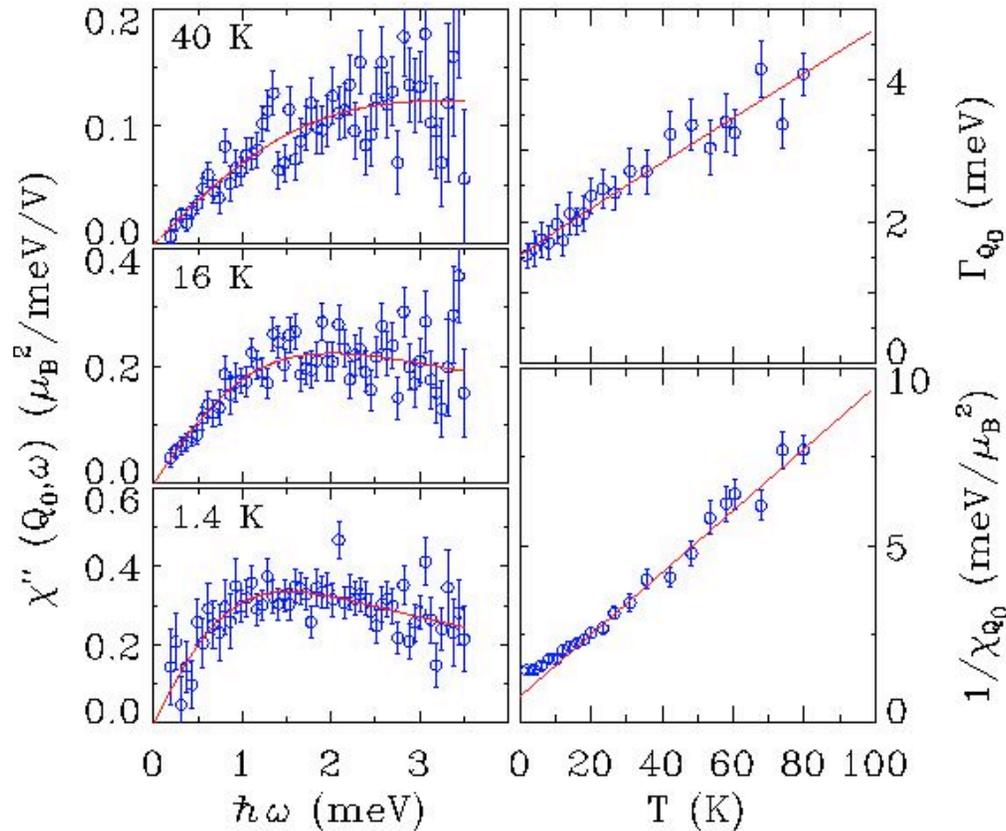
LiV₂O₄ (3d^{1.5}): Dynamic Spin Correlations



SHL et al., PRL (2001)

**Dynamic spin correlations become AFM
as LiV₂O₄ enters the heavy fermion phase**

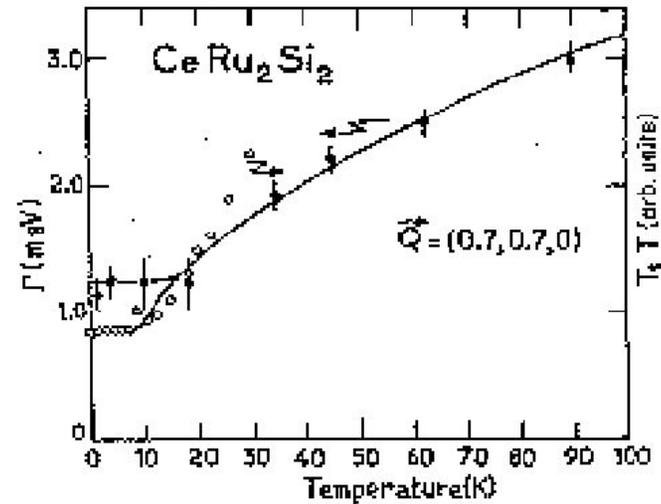
Dynamic Magnetic Correlations in LiV_2O_4



$$\Gamma(T) = \Gamma_0 + k_B T^\alpha$$

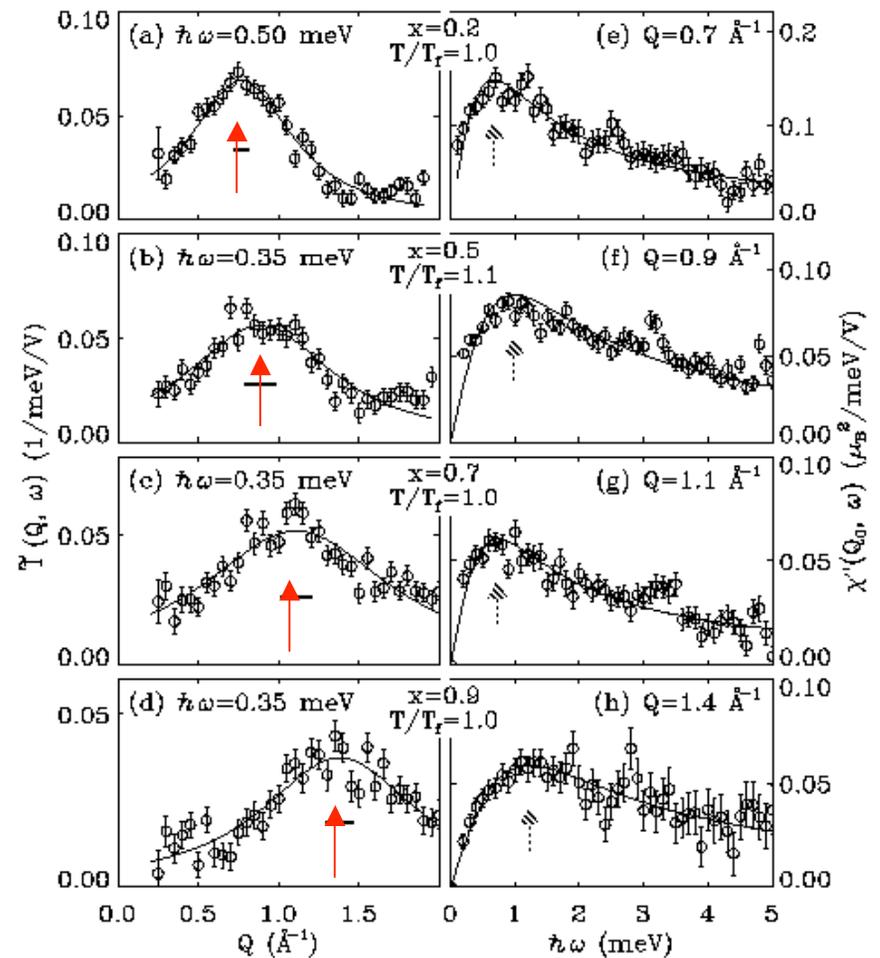
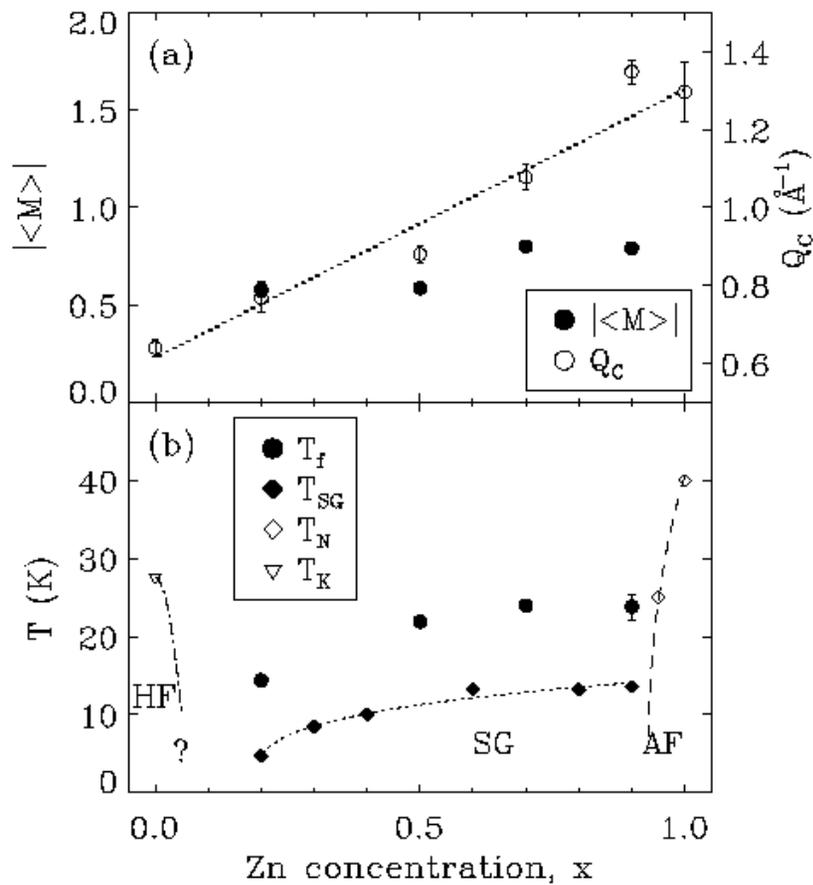
with $\Gamma_0 = 1.5 \text{ meV} = T_K$ (Kondo T) $\sim T_c$
 and $\alpha = 0.96(2)$

L.P. Regnault et al., PRB 38, 4481 (1988)



Linear $\Gamma(T)$ is a common feature of geometrically frustrated antiferromagnets.

Evolution of magnetic correlations in $\text{Li}_{1-x}\text{Zn}_x\text{V}_2\text{O}_4$



S. Park et al., unpublished (2003)

Summary



1. Ground state degeneracy
→ Spectral weight down-shift to $E < \hbar\omega_{\text{cw}}$
2. Weak connectivity of spin correlations due to triangular motif
→ Emergence of local spin clusters with zero energy modes



3. Relation between the metallic properties and magnetism is still in question.
The clue to the answer might be in understanding why unlike insulating magnets the characteristic wave vector for magnetic correlations changes gradually as the Zn concentration

Elastic and Inelastic diffuse scattering on single crystals is indispensable to identify the fundamental spin degrees of freedom in GF magnets